# Intro to High Spectral Resolution IR Measurements 

Lectures in Madison<br>27 March 2013

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line broadening with pressure helps to explain weighting functions

line broadening with pressure helps to explain weighting functions

$$
-\mathrm{k}_{\mathrm{v}} \mathrm{u}(\mathrm{z})
$$

$$
\tau_{v}(\mathrm{z} \rightarrow \infty)=\mathrm{e}
$$



Wavenumber
Energy Contribution

For a given water vapor spectral channel the weighting function depends on the amount of water vapor in the atmospheric column


CO 2 is about the same everywhere, the weighting function for a given CO 2 spectral channel is the same everywhere

## Vibrational Bands



## $\mathrm{CO}_{2}$ Vibration - Rotation Spectra

$$
E(v, J)=\underbrace{h v\left(v+\frac{1}{2}\right)-x h v\left(v+\frac{1}{2}\right)^{2}+\ldots}_{\text {vibration }}+\underbrace{B_{v}\left[J(J+1)-\ell^{2}\right]-D_{v}\left[J(J+1)-\ell^{2}\right]^{2}+\ldots}_{\text {rotation }}
$$



## $\mathrm{H}_{2} \mathrm{O}$ Vibration - Rotation Spectra




## Rotational Lines



## Earth emitted spectrum in CO2 sensitive 705 to 760 cm-1



## Broad Band


... in Brightness Temperature


## High Spectral Resolution



Sampling over rotational bands

## Infrared Radiance and Brightness Temperature Spectrum

Planck Function

$$
B_{v}(T)=2 h c^{2} v^{3} / \exp \left(\frac{h c v}{k T}\right)-1
$$

Upwelling IR radiation

$$
R_{v}=\int_{z_{0}}^{\infty} B_{v}(T(z)) \frac{d \tau_{v}(z)}{d z} d z
$$




From E. Weisz

## Atmospheric Temperafure Profile Retrieval

$$
R_{v}=\int_{p s}^{0} B_{v}(T(p)) W_{v}(p) d p
$$

$$
W_{v}(p)=\partial \tau_{v}(p) / \partial \ln p
$$



High-spectral measurements



Profiles at high-vertical resolution
From E. Weisz

## Regression Retrieval Summary



From E. Weisz

## Dual-Regression Retrieval





Moisture Weighting Functions
High spectral resolution advanced sounder will have more and sharper weighting functions compared to current GOES sounder. Retrievals will have better vertical resolution.


## Resolving absorption features in atmospheric windows

 enables detection of temperature inversions

Detection of inversions is critical for severe weather forecasting. Combined with improved low-level moisture depiction, key ingredients for night-time severe storm development can be monitored.



## Ability to detect inversions disappears with broadband observations (> $3 \mathrm{~cm}-1$ )



Longwave window region


Longwave window region


Longwave window region


Longwave window region


Longwave window region


Longwave window region


Longwave window region

## Twisted Ribbon formed by $\mathrm{CO}_{2}$ spectrum:

 Tropopause inversion causes On-line \& off-line patterns to cross

Wavelength ( $\mu \mathrm{m}$ )



Inferring surface properties with AIRS high spectral resolution data Barren region detection if T1086 < T981

## Barren vs Water/Vegetated



AIRS data from 14 June 2002
$\mathrm{T}\left(981 \mathrm{~cm}^{-1}\right)-\mathrm{T}\left(1086 \mathrm{~cm}^{-1}\right)$

$\mathrm{T}\left(1086 \mathrm{~cm}^{-1}\right)$

from Tobin et al.

$R=\varepsilon_{s} B s(1-\sigma c)+\sigma_{c} B c \quad$ using $e^{-\sigma}=1-\sigma$

So difference of thin ice cloud over ocean from clear sky over ocean is given by
$\Delta \mathrm{R}=-\varepsilon s \sigma_{c} \mathrm{~B} s+\sigma c \mathrm{Bc}$

For $\mathrm{Bs}>\mathrm{Bc}$ and $\varepsilon s \sim 1$
$\Delta R=-\sigma_{c} B s+\sigma_{c} B c=\sigma_{c}[B c-B s]$

As $\boldsymbol{\sigma c}$ increases (decreases) then $\Delta \mathrm{R}$ becomes more negative (positive)


Imaginary Index of Refraction of Ice and Dust


## IASI detection of dust <br> 

red spectrum is from nearby clear fov

IASI detection of cirrus


wavenumber $1349.75 \mathrm{~cm}^{\mathrm{A}}$-1
 window

- In the 800-1000 $\mathrm{cm}^{-1}$
atmospheric window:
Silicate index increases
Ice index decreases
with wavenumber

Volz, F.E. : Infrared optical constant of
ammonium sulphate, Sahara Dust, ammonium sulphate, Sahara Dust,
volcanic pumice and flash, Appl Opt 12
$564-658$ (1973)

AIRS.2002.10.28.123.L1B.AIRS_Rad.v2.6.10.3.A02302200913
$\sim 12521 / \mathrm{cm} \mathrm{Tb}-\sim 9131 / \mathrm{cm} \mathrm{Tb}$


AIRS.2002.10.28.123.L1B.AIRS_Rad.v2.6.10.3.A02302200913
$\sim 9131 / \mathrm{cm} \mathrm{Tb}-\sim 8371 / \mathrm{cm} \mathrm{Tb}$



## Mt Etna Ash cloud at 500 hPa


wavenumber $919.47 \mathbf{c m}^{\wedge}$ - 1

istrument: AIRS






# Mt Etna volcanic plume SO2 (left) from 1284-1345 <br> Ash (right) from 832-900 



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## AIRS Spectra from around the Globe



# Intro to Microwave and Split Window Moisture 

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## Earth emitted spectra overlaid on Planck function envelopes

High resolution atmospheric absorption spectrum and comparative blackbody curves.


## MODIS IR Spectral Bands

High resolution atmospheric absorption spectrum and comparative blackbody curves.


## First order estimation of SST correcting for low level moisture

Moisture attenuation in atmospheric windows varies linearly with optical depth.

$$
\tau_{\lambda}=\mathrm{e}^{-\mathrm{k}_{\lambda} \mathrm{u}}=1-\mathrm{k}_{\lambda} \mathrm{u}
$$

For same atmosphere, deviation of brightness temperature from surface temperature is a linear function of absorbing power. Thus moisture corrected SST can inferred by using split window measurements and extrapolating to zero $\mathrm{k}_{\lambda}$

$$
\mathrm{T}_{\mathrm{s}}=\mathrm{T}_{\mathrm{bw} 1}+\left[\mathrm{k}_{\mathrm{w} 1} /\left(\mathrm{k}_{\mathrm{w} 2}-\mathrm{k}_{\mathrm{w} 1}\right)\right]\left[\mathrm{T}_{\mathrm{bw} 1}-\mathrm{T}_{\mathrm{bw} 2}\right]
$$

Moisture content of atmosphere inferred from slope of linear relation.


## MODIS SEA SURFACE TEMPERATURE




## In the IRW - A is off H 2 O line and B is on H 2 O line




Weighting Function

## Radiation is governed by Planck's Law

$$
\mathbf{B}(\lambda, T)=c_{1} /\left\{\lambda^{5}\left[e^{c_{2} / \lambda T}-1\right]\right\}
$$

In microwave region $c_{2} / \lambda T \ll 1$ so that

$$
\mathbf{e}^{\mathbf{c}_{2} / \lambda T}=1+\mathbf{c}_{2} / \lambda T+\text { second order }
$$

And classical Rayleigh Jeans radiation equation emerges

$$
\mathbf{B}_{\lambda}(\mathbf{T}) \approx\left[\mathbf{c}_{1} / \mathbf{c}_{2}\right]\left[\mathbf{T} / \lambda^{4}\right]
$$

Radiance is linear function of brightness temperature.

## ISCCP-Dx 19g207-199sigi Hean Annual



## ISCCP-D1 1992 Hean Annual



10 to 11 um

## Microwave Form of RTE

$$
\begin{array}{r}
\begin{array}{r}
\mathrm{I}^{\mathrm{sfc}}=\varepsilon_{\lambda} \mathrm{B}_{\lambda}\left(\mathrm{T}_{\mathrm{s}}\right) \tau_{\lambda}\left(\mathrm{p}_{\mathrm{s}}\right) \\
\lambda
\end{array}+\left(1-\varepsilon_{\lambda}\right) \tau_{\lambda}\left(\mathrm{p}_{\mathrm{s}}\right) \int_{\mathrm{o}}^{\mathrm{p}_{\mathrm{s}}} \mathrm{~B}_{\lambda}(\mathrm{T}(\mathrm{p}))
\end{array} \frac{\partial \tau_{\lambda}^{\prime}(\mathrm{p})}{\partial \ln \mathrm{p}} \mathrm{~d} \ln \mathrm{p} .
$$



In the microwave region $c_{2} / \lambda T \ll 1$, so the Planck radiance is linearly proportional to the brightness temperature

$$
\mathrm{B}_{\lambda}(\mathrm{T}) \approx\left[\mathrm{c}_{1} / \mathrm{c}_{2}\right]\left[\mathrm{T} / \lambda^{4}\right]
$$

So

$$
\mathrm{T}_{\mathrm{b} \lambda}=\varepsilon_{\lambda} \mathrm{T}_{\mathrm{s}}\left(\mathrm{p}_{\mathrm{s}}\right) \tau_{\lambda}\left(\mathrm{p}_{\mathrm{s}}\right)+\int_{\mathrm{p}_{\mathrm{s}}}^{\mathrm{o}} \mathrm{~T}(\mathrm{p}) \mathrm{F}_{\lambda}(\mathrm{p}) \frac{\partial \tau_{\lambda}(\mathrm{p})}{\partial \ln \mathrm{p}} \mathrm{~d} \ln \mathrm{p}
$$

where

$$
\mathrm{F}_{\lambda}(\mathrm{p})=\left\{1+\left(1-\varepsilon_{\lambda}\right)\left[\frac{\tau_{\lambda}\left(\mathrm{p}_{\mathrm{s}}\right)}{\tau_{\lambda}(\mathrm{p})}\right]^{2}\right\}
$$

## Transmittance

$$
\begin{aligned}
& \tau(\mathrm{a}, \mathrm{~b})=\tau(\mathrm{b}, \mathrm{a}) \\
& \tau(\mathrm{a}, \mathrm{c})=\tau(\mathrm{a}, \mathrm{~b}) * \tau(\mathrm{~b}, \mathrm{c})
\end{aligned}
$$

Thus downwelling in terms of upwelling can be written
$\tau^{\prime}(\mathrm{p}, \mathrm{ps})=\tau(\mathrm{ps}, \mathrm{p})=\tau(\mathrm{ps}, 0) / \tau(\mathrm{p}, 0)$
and

$$
\mathrm{d} \tau^{\prime}(\mathrm{p}, \mathrm{ps})=-\mathrm{d} \tau(\mathrm{p}, 0) * \tau(\mathrm{ps}, 0) /[\tau(\mathrm{p}, 0)]^{2}
$$

| WAVELENGTH |  | FREQUENCY |  | WAVENUMBER |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{cm} \quad \mu \mathrm{m}$ | Å | Hz | GHz | $\mathrm{cm}^{-1}$ |
| $10^{-5}$ Near Ultraviolet (UV) | 1,000 | $3 \times 10^{15}$ |  |  |
| $4 \times 10^{-5}$ 0.4 <br> Visible  | 4,000 | $7.5 \times 10^{14}$ |  |  |
| $7.5 \times 10^{-5}$ Near Infrared (IR) | 7,500 | $4 \times 10^{14}$ |  | 13,333 |
| $\begin{array}{ll} 2 \times 10^{-3} \\ \text { Far Infrared (IR) } \end{array}$ | $2 \times 10^{5}$ | $1.5 \times 10^{13}$ |  | 500 |
| $\begin{array}{cc} 0.1 & 10^{3} \\ \text { Microwave (MW) } \end{array}$ |  | $3 \times 10^{11}$ | 300 | 10 |




## Microwave spectral bands

23.8 GHz dirty window H2O absorption
31.4 GHz window
$60 \mathrm{GHz} \quad \mathrm{O} 2$ sounding
$120 \mathrm{GHz} \quad \mathrm{O} 2$ sounding
$183 \mathrm{GHz} \quad \mathrm{H} 2 \mathrm{O}$ sounding

$23.8,31.4,50.3,52.8,53.6,54.4,54.9,55.5,57.3$ ( 6 chs ), $89.0 \stackrel{57}{\mathrm{GH}}$



$\mathrm{Tb}=\varepsilon_{\mathrm{s}} \mathrm{T}_{\mathrm{s}} \tau_{\mathrm{m}}+\varepsilon_{\mathrm{m}} \mathrm{T}_{\mathrm{m}}+\varepsilon_{\mathrm{m}} \mathbf{r}_{\mathrm{s}} \tau_{\mathrm{m}} \mathrm{T}_{\mathrm{m}}$
$\mathrm{Tb}=\boldsymbol{\varepsilon}_{\mathrm{s}} \mathrm{T}_{\mathrm{s}}(1-\sigma \mathrm{m})+\boldsymbol{\sigma}_{\mathrm{m}} \mathrm{T}_{\mathrm{m}}+\boldsymbol{\sigma m}\left(1-\varepsilon_{\mathrm{s}}\right)(1-\boldsymbol{\sigma m}) \mathrm{Tm} \quad$ using $\mathrm{e}^{-\sigma}=1-\sigma$
So temperature difference of low moist over ocean from clear sky over ocean is given by
$\Delta \mathrm{Tb}=-\boldsymbol{\varepsilon} s \boldsymbol{\sigma}_{\mathrm{m}} \mathrm{T}_{\mathrm{s}}+\boldsymbol{\sigma} \mathrm{m} \mathrm{T}_{\mathrm{m}}+\boldsymbol{\sigma} \mathrm{m}(1-\boldsymbol{\varepsilon} \mathrm{s})\left(1-\boldsymbol{\sigma}_{\mathrm{m}}\right) \mathrm{T}_{\mathrm{m}}$
For $\varepsilon_{s} \sim 0.5$ and $\mathrm{T}_{\mathrm{s}} \sim \mathrm{T}_{\mathrm{m}}$ this is always positive for $0<\sigma_{\mathrm{m}}<1$

$\mathrm{R}=\varepsilon_{s} \mathrm{Bs}(1-\sigma \mathrm{m})+\sigma \mathrm{m} \mathrm{Bm} \quad$ using $\mathrm{e}^{-\sigma}=1-\sigma$ and $\tau \sim 1-\sigma \sim 1-\mathrm{a}$

So difference of low mist over ocean from clear sky over ocean is given by
$\Delta \mathrm{R}=-\boldsymbol{\varepsilon} s \sigma_{\mathrm{m}} \mathrm{Bs}+\sigma \mathrm{m} B m$

For $\boldsymbol{\varepsilon} s \sim 1$
$\Delta \mathrm{R}=-\boldsymbol{\sigma} \mathrm{m} \mathrm{s}+\boldsymbol{\sigma} \mathrm{m} \mathrm{Bm}=\boldsymbol{\sigma} \mathrm{m}[\mathrm{Bm}-\mathrm{B} s]$

So if $[\mathrm{Bm}-\mathrm{Bs}]<0$ then as $\sigma \mathrm{m}$ increases $\Delta \mathrm{R}$ becomes more negative






## ATMS Weighting Functions





Spectral regions used for remote sensing of the earth atmosphere and surface from satellites. $\varepsilon$ indicates emissivity, q denotes water vapour, and T represents temperature.

