

Investigations with Remote Sensing Data using MODIS, AIRS, & AMSU (4 Labs)



Paul Menzel UW and colleagues at CIMSS Access to text, visualization tools, and data

For Remote Sensing Applications Text ftp://ftp.ssec.wisc.edu/pub/menzel/AppMetSat09.pdf

For HYDRA http://www.ssec.wisc.edu/hydra/

For MODIS data and quick browse images http://rapidfire.sci.gsfc.nasa.gov/realtime

For MODIS data http://ladsweb.nascom.nasa.gov/

For AIRS and AMSU data http://daac.gsfc.nasa.gov/

HYperspectral viewer for Development of Research Applications - HYDRA MODI

MSG, GOES

Freely available software For researchers and educators Computer platform independent Extendable to more sensors and applications Based in VisAD (Visualization for Algorithm Development) Uses Jython (Java implementation of Python) runs on most machines 512MB main memory & 32MB graphics card suggested on-going development effort

Rink et al, BAMS 2007



MODIS, AIRS, IASI, AMSU, CALIPSO

Developed at CIMSS by Tom Rink Tom Whittaker Kevin Baggett

With guidance from Paolo Antonelli Liam Gumley Paul Menzel Allen Huang



http://www.ssec.wisc.edu/hydra/

Intro to VIS-IR Radiation Lab

Lecture - Labs in Italy 18 – 25 May 2012 in Bologna 28 May - 1 June 2012 in Potenza

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Relevant Material in Applications of Meteorological Satellites

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Satellite remote sensing of the Earth-atmosphere



Observations depend on

telescope characteristics (resolving power, diffraction) detector characteristics (field of view, signal to noise) communications bandwidth (bit depth) spectral intervals (window, absorption band) time of day (daylight visible) atmospheric state (T, Q, clouds) earth surface (Ts, vegetation cover)

Electromagnetic spectrum



Spectral Characteristics of Energy Sources and Sensing Systems



Definitions of Radiation

QUANTITY	SYMBOL	UNITS
Energy	dQ	Joules
Flux	dQ/dt	Joules/sec = Watts
Irradiance	dQ/dt/dA	Watts/meter ²
Monochromatic Irradiance	dQ/dt/dA/dλ	W/m ² /micron
11 i uululiee	or	
	dQ/dt/dA/dv	W/m ² /cm ⁻¹
Radiance	$dQ/dt/dA/d\lambda/d\Omega$	W/m ² /micron/ster
	or	
	dQ/dt/dA/dv/dΩ	W/m²/cm ⁻¹ /ster

Using wavelengths

 $c_2/\lambda T$

Planck's Law

 $B(\lambda,T) = c_1 / \lambda^5 / [e -1] \quad (mW/m^2/ster/cm)$

where

$$\lambda = \text{ wavelengths in cm}$$

$$T = \text{temperature of emitting surface (deg K)}$$

$$c_1 = 1.191044 \text{ x } 10-5 \text{ (mW/m²/ster/cm-4)}$$

$$c_2 = 1.438769 \text{ (cm deg K)}$$

Wien's Law $dB(\lambda_{max},T) / d\lambda = 0$ where $\lambda(max) = .2897/T$ indicates peak of Planck function curve shifts to shorter wavelengths (greater wavenumbers)with temperature increase. Note $B(\lambda_{max},T) \sim T^5$.

Stefan-Boltzmann Law $E = \pi \int B(\lambda,T) d\lambda = \sigma T^4$, where $\sigma = 5.67 \times 10^{-8} \text{ W/m}^2/\text{deg}^4$.

states that irradiance of a black body (area under Planck curve) is proportional to T^4 .

Brightness Temperature

$$T = c_2 / \left[\lambda \ln(\frac{c_1}{-+} + 1) \right]$$
 is determined by inverting Planck function
$$\frac{\lambda^5 B_{\lambda}}{10}$$

Spectral Distribution of Energy Radiated from Blackbodies at Various Temperatures



Planck Tool







Using wavenumbers

Planck's Law where $B(v,T) = c_1v^3/[e -1]$ (mW/m²/ster/cm⁻¹) v = # wavelengths in one centimeter (cm-1) T = temperature of emitting surface (deg K) $c_1 = 1.191044 \times 10-5$ (mW/m²/ster/cm⁻⁴) $c_2 = 1.438769$ (cm deg K)

Wien's Law $dB(v_{max},T) / dv = 0$ where v_{max}) = 1.95T

indicates peak of Planck function curve shifts to shorter wavelengths (greater wavenumbers) with temperature increase.

Stefan-Boltzmann Law $E = \pi \int_{0}^{\infty} B(v,T) dv = \sigma T^4$, where $\sigma = 5.67 \times 10^{-8} \text{ W/m}^2/\text{deg}^4$.

states that irradiance of a black body (area under Planck curve) is proportional to T^4 .

Brightness Temperature

$$\Gamma = c_2 v / [ln(---+1)] \text{ is determined by inverting Planck function} B_v$$

Using wavenumbers

$$c_2 v/T$$

B(v,T) = c_1 v³ / [e -1]
(mW/m²/ster/cm⁻¹)

v(max in cm-1) = 1.95T

 $B(v_{max},T) \sim T^{**3}$.

$$E = \pi \int B(v,T) dv = \sigma T^{4},$$

$$O = \frac{c_{1}v^{3}}{T = c_{2}v/[\ln(-+1)]}$$

$$B_{v}$$



Using wavelengths

 $c_{2}/\lambda T$ $B(\lambda,T) = c_{1}/\{ \lambda^{5} [e -1] \}$ $(mW/m^{2}/ster/\mu m)$

 $\lambda(\max \text{ in cm})T = 0.2897$

B(λ_{max} ,T) ~ T**5.



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Planck Tool





Temperature Sensitivity of $B(\lambda,T)$ for typical earth temperatures





(Approximation of) B as function of α and T

$\Delta B/B = \alpha \Delta T/T$

Integrating the Temperature Sensitivity Equation Between T_{ref} and T (B_{ref} and B):

 $B=B_{ref}(T/T_{ref})^{\alpha}$

Where $\alpha = c_2 v / T_{ref}$ (in wavenumber space)



The temperature sensitivity indicates the power to which the Planck radiance depends on temperature, since B proportional to T^{α} satisfies the equation. For infrared wavelengths,

 $\alpha = c_2 \nu / T = c_2 / \lambda T.$

Wavenumber	Typical Scene Temperature	Temperature Sensitivity
900	300	4.32
2500	300	11.99



Non-Homogeneous FOV

For NON-UNIFORM FOVs:

 $B_{obs} = NB_{cold} + (1-N)B_{hot}$

$$B_{obs} = N B_{ref} (T_{cold}/T_{ref})^{\alpha} + (1-N) B_{ref} (T_{hot}/T_{ref})^{\alpha}$$



$$B_{obs} = B_{ref} (1/T_{ref})^{\alpha} (N T_{cold}^{\alpha} + (1-N)T_{hot}^{\alpha})$$

For N=.5

$$B_{obs}/B_{ref} = .5 (1/T_{ref})^{\alpha} (T_{cold}^{\alpha} + T_{hot}^{\alpha})$$

 $B_{obs}/B_{ref} = .5 (1/T_{ref}T_{cold})^{\alpha} (1 + (T_{hot}/T_{cold})^{\alpha})$

The greater α the more predominant the hot term

At 4 μ m (α =12) the hot term more dominating than at 11 μ m (α =4)



Cloud edges and broken clouds appear different in 11 and 4 um images.

 $T(11)^{**}4 = (1-N)^{*}Tclr^{**}4 + N^{*}Tcld^{**}4 \sim (1-N)^{*}300^{**}4 + N^{*}200^{**}4$ $T(4)^{**}12 = (1-N)^{*}Tclr^{**}12 + N^{*}Tcld^{**}12 \sim (1-N)^{*}300^{**}12 + N^{*}200^{**}12$

Cold part of pixel has more influence for B(11) than B(4)

Relevant Material in Applications of Meteorological Satellites

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Solar (visible) and Earth emitted (infrared) energy



Incoming solar radiation (mostly visible) drives the earth-atmosphere (which emits infrared).

Over the annual cycle, the incoming solar energy that makes it to the earth surface (about 50 %) is balanced by the outgoing thermal infrared energy emitted through the atmosphere.

The atmosphere transmits, absorbs (by H2O, O2, O3, dust) reflects (by clouds), and scatters (by aerosols) incoming visible; the earth surface absorbs and reflects the transmitted visible. Atmospheric H2O, CO2, and O3 selectively transmit or absorb the outgoing infrared radiation. The outgoing microwave is primarily affected by H2O and O2.

Spectral Characteristics of Atmospheric Transmission and Sensing Systems







Normalized black body spectra representative of the sun (left) and earth (right), plotted on a logarithmic wavelength scale. The ordinate is multiplied by wavelength so that the area under the curves is proportional to irradiance. 27





BT11=290K and BT4=310K. What fraction of R4 is due to reflected solar radiance?

```
R4 = R4 \text{ refl} + R4 \text{ emiss}BT4 \text{ emiss} = BT11R4 \sim T^{**}12
```

Fraction = $[310^{**}12 - 290^{**}12]/310^{**}12 \sim .55$



Infrared (Emissive Bands)

Radiative Transfer Equation in the IR



Emission, Absorption, Reflection, and Scattering

Blackbody radiation B_{λ} represents the upper limit to the amount of radiation that a real substance may emit at a given temperature for a given wavelength.

Emissivity ϵ_{λ} is defined as the fraction of emitted radiation R_{λ} to Blackbody radiation,

$$\varepsilon_{\lambda} = R_{\lambda} / B_{\lambda}$$

In a medium at thermal equilibrium, what is absorbed is emitted (what goes in comes out) so

 $a_{\lambda} = \varepsilon_{\lambda}$.

Thus, materials which are strong absorbers at a given wavelength are also strong emitters at that wavelength; similarly weak absorbers are weak emitters.

If a_{λ} , r_{λ} , and τ_{λ} represent the fractional absorption, reflectance, and transmittance, respectively, then conservation of energy says

$$a_{\lambda} + r_{\lambda} + \tau_{\lambda} = 1$$
 .

For a blackbody $a_{\lambda} = 1$, it follows that $r_{\lambda} = 0$ and $\tau_{\lambda} = 0$ for blackbody radiation. Also, for a perfect window $\tau_{\lambda} = 1$, $a_{\lambda} = 0$ and $r_{\lambda} = 0$. For any opaque surface $\tau_{\lambda} = 0$, so radiation is either absorbed or reflected $a_{\lambda} + r_{\lambda} = 1$.

At any wavelength, strong reflectors are weak absorbers (i.e., snow at visible wavelengths), and weak reflectors are strong absorbers (i.e., asphalt at visible wavelengths). 32

 $\mathbf{r}_{\!\lambda}\mathbf{R}_{\!\lambda}$

 $\tau_{\lambda} R_{\lambda}$

R

 $\epsilon_{\lambda}\mathsf{B}_{\lambda}(\mathsf{T})$

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Transmittance

Transmission through an absorbing medium for a given wavelength is governed by the number of intervening absorbing molecules (path length u) and their absorbing power (k_{λ}) at that wavelength. Beer's law indicates that transmittance decays exponentially with increasing path length

$$\tau_{\lambda} (z \to \infty) = e^{-k_{\lambda} u(z)}$$

where the path length is given by $u(z) = \int_{-\infty}^{\infty} \rho dz$.

 k_{λ} u is a measure of the cumulative depletion that the beam of radiation has experienced as a result of its passage through the layer and is often called the optical depth σ_{λ} .

Z

Realizing that the hydrostatic equation implies $g \rho dz = -q dp$

where q is the mixing ratio and ρ is the density of the atmosphere, then

$$\mathfrak{u}(p) = \int_{0}^{p} q g^{-1} dp \quad \text{and} \quad \tau_{\lambda}(p \to o) = e^{-k_{\lambda} u(p)}$$

Spectral Characteristics of Atmospheric Transmission and Sensing Systems




Aerosol Size Distribution

There are 3 modes :

- « nucleation »: radius is between 0.002 and 0.05 μ m. They result from combustion processes, photo-chemical reactions, etc.

- « accumulation »: radius is between 0.05 μm and 0.5 μm. Coagulation processes.

- « **coarse** »: larger than 1 μm. From mechanical processes like aeolian erosion.

« fine » particles (nucleation and accumulation) result from anthropogenic activities, coarse particles come from natural processes.



Measurements in the Solar Reflected Spectrum across the region covered by AVIRIS



AVIRIS Movie #1

AVIRIS Image - Linden CA 20-Aug-1992 224 Spectral Bands: 0.4 - 2.5 μm Pixel: 20m x 20m Scene: 10km x 10km





AVIRIS Movie #2

AVIRIS Image - Porto Nacional, Brazil 20-Aug-1995 224 Spectral Bands: 0.4 - 2.5 μm Pixel: 20m x 20m Scene: 10km x 10km





Intro to Land-Ocean-Atmosphere Remote Sensing Lab

Lecture - Labs in Bologna & Potenza May-June 2012

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Re-emission of Infrared Radiation



Molecular Responses to Radiation



Molecular absorption of IR by vibrational and rotational excitation



Earth emitted spectra overlaid on Planck function envelopes



Earth emitted spectra overlaid on Planck function envelopes



High resolution atmospheric absorption spectrum and comparative blackbody curves.

Radiative Transfer Equation

The radiance leaving the earth-atmosphere system sensed by a satellite borne radiometer is the sum of radiation emissions from the earth-surface and each atmospheric level that are transmitted to the top of the atmosphere. Considering the earth's surface to be a blackbody emitter (emissivity equal to unity), the upwelling radiance intensity, I_{λ} , for a cloudless atmosphere is given by the expression

$$\begin{split} I_{\lambda} &= \epsilon_{\lambda}{}^{sfc} \ B_{\lambda}(\ T_{sfc}) \ \tau_{\lambda}(sfc \ - \ top) \ + \ \sum \epsilon_{\lambda}{}^{layer} \ B_{\lambda}(\ T_{layer}) \ \tau_{\lambda}(layer \ - \ top) \\ layers \end{split}$$

where the first term is the surface contribution and the second term is the atmospheric contribution to the radiance to space.

Satellite observation comes from the sfc and the layers in the atm



esfc for earth surface

recalling that $\varepsilon i = 1$ - τi for each layer, then

Robs = ε sfc Bsfc $\tau 1 \tau 2 \tau 3 + (1-\tau 1) B1 \tau 2 \tau 3 + (1-\tau 2) B2 \tau 3 + (1-\tau 3) B3$

 $I_{\lambda} = \varepsilon_{\lambda}^{sfc} B_{\lambda}(T_{sfc}) \tau_{\lambda}(sfc - top) + \sum \varepsilon_{\lambda}^{layer} B_{\lambda}(T_{layer}) \tau_{\lambda}(layer - top)$ layers

The emission of an infinitesimal layer of the atmosphere at pressure p is equal to the absorption (1 - transmission). So,

 $\varepsilon_{\lambda}(\text{layer}) \tau_{\lambda}(\text{layer to top}) = [1 - \tau_{\lambda}(\text{layer})] \tau_{\lambda}(\text{layer to top})$

Since transmission is multiplicative $\tau_{\lambda}(\text{layer to top}) - \tau_{\lambda}(\text{layer}) \tau_{\lambda}(\text{layer to top}) = -\Delta \tau_{\lambda}(\text{layer to top})$

So we can write

$$\begin{split} I_{\lambda} &= \epsilon_{\lambda}{}^{sfc} B_{\lambda}(T(p_{s})) \tau_{\lambda}(p_{s}) - \Sigma B_{\lambda}(T(p)) \Delta \tau_{\lambda}(p) \ . \\ p \\ \text{which when written in integral form reads} \\ I_{\lambda} &= \epsilon_{\lambda}{}^{sfc} B_{\lambda}(T(p_{s})) \tau_{\lambda}(p_{s}) - \int_{0}^{p_{s}} B_{\lambda}(T(p)) \left[d\tau_{\lambda}(p) / dp \right] dp \ . \\ q \\ 0 \end{split}$$

Weighting Functions



Weighting Functions





CO2 channels see to different levels in the atmosphere



14.2 um 13.9 um 13.6 um 13.3 um

Characteristics of RTE

- * Radiance arises from deep and overlapping layers
- * The radiance observations are not independent
- There is no unique relation between the spectrum of the outgoing radiance and T(p) or Q(p)
- * T(p) is buried in an exponent in the denominator in the integral
- * Q(p) is implicit in the transmittance
- Boundary conditions are necessary for a solution; the better the first guess the better the final solution

Profile Retrieval from Sounder Radiances

$$I_{\lambda} = \epsilon_{\lambda}^{sfc} B_{\lambda}(T(p_s)) \tau_{\lambda}(p_s) - \int_{0}^{p_s} B_{\lambda}(T(p)) F_{\lambda}(p) \left[\frac{d\tau_{\lambda}(p)}{dp} \right] dp.$$

I1, I2, I3,, In are measured with the sounding radiometer P(sfc) and T(sfc) come from ground based conventional observations $\tau_{\lambda}(p)$ are calculated with physics models (using for CO2 and O3) $\varepsilon_{\lambda}^{sfc}$ is estimated from a priori information (or regression guess)

First guess solution is inferred from (1) in situ radiosonde reports, (2) model prediction, or (3) blending of (1) and (2)

Profile retrieval from perturbing guess to match measured sounder radiances

Example Sounding



Viewing remote sensing data with HYDRA





Relevant Material in Applications of Meteorological Satellites

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High clouds reflect more than surface at 0.65 μ m



High clouds, cooler than surface, create lower 11 µm BTs



⁶¹ High clouds and snow both reflect a lot at 0.65 μ m



⁶² High clouds reflect but snow doesn't at 1.64 μ m





[†] Low clouds, cooler than surface, create lower 11 µm BTs



Low clouds reflecting create larger 4 µm brightness temperatures





Detecting low clouds in 4-11 µm brightness temperature differences



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Detecting ice clouds in 8.6-11 µm brightness temperature differences



Optical properties of cloud particles: imaginary part of refraction index

Imaginary part of refraction index



⁷⁰ BT[8.6] - BT[11] will be positive for transmissive ice clouds



MODIS identifies cloud classes



Hi cld Mid cld Lo cld Snow
Clouds separate into classes when multispectral radiance information is viewed



Cloud Mask Tests

- BT11
- BT13.9
- BT6.7
- BT3.9-BT11
- BT11-BT12
- BT8.6-BT11
- BT6.7-BT11 or BT13.9-BT11
- BT11+aPW(BT11-BT12)
- r0.65
- r0.85
- r1.38
- r1.6
- r0.85/r0.65 or NDVI
- σ(BT11)

clouds over ocean high clouds high clouds broken or scattered clouds high clouds in tropics ice clouds clouds in polar regions clouds over ocean clouds over land clouds over ocean thin cirrus clouds over snow, ice cloud clouds over vegetation clouds over ocean

ATMOSPHERE - THERMAL RADIATION







ATMOSPHERE-SOLAR RADIATION



C351.008 5/93



MAS (SUCCESS) 1996/04/26 18:43:48 UTC Track 03, Band 45 (11.01 micron) Brightness Temp. (K)





MAS (SUCCESS) 1996/04/26 18:43:48 UTC Track 03, Band 15 (1.90 micron) Reflectance





OCEAN-SOLAR RADIATION



MODIS SEA SURFACE TEMPERATURE



EOS



The warm heart of the Gulf Stream is readily apparent in the top SST image. As the current flows toward the northeast it begins to meander and pinch off eddies that transport warm water northward and cold water southward. The current also divides the local ocean into a low-biomass region to the south and a higher-biomass region to the north. The data were collected by MODIS aboard Aqua on April 18, 2005.

LAND - THERMAL RADIATION



EOS



LAND-SOLAR RADIATION



WAVELENGTH (λ) NANOMETERS

Example with MODIS



WAVELENGTH (λ) NANOMETERS 83

2500





Investigating with Multi-spectral Combinations

Given the spectral response of a surface or atmospheric feature

Select a part of the spectrum where the reflectance or absorption changes with wavelength

e.g. reflection from grass

If 0.65 μm and 0.85 μm channels see the same reflectance than surface viewed is not grass; if 0.85 μm sees considerably higher reflectance than 0.65 μm then surface might be grass

Seasonal Biosphere Ocean Chlorophyll-a & Terrestrial NDVI



High resolution atmospheric absorption spectrum and comparative blackbody curves.



Intro to Lab on High Spectral Resolution IR Measurements

Lecture - Labs in Bologna & Potenza May-June 2012

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line broadening with pressure helps to explain weighting functions



line broadening with pressure helps to explain weighting functions



For a given water vapor spectral channel the weighting function depends on the amount of water vapor in the atmospheric column



CO2 is about the same everywhere, the weighting function for a given CO2 spectral channel is the same everywhere ⁹⁰

Vibrational Bands





D. Tobin, UMBC





Rotational Lines



Earth emitted spectrum in CO2 sensitive 705 to 760 cm-1



Broad Band





High Spectral Resolution



Sampling over rotational bands



Moisture Weighting Functions

High spectral resolution advanced sounder will have more and sharper weighting functions compared to current GOES sounder. Retrievals will have better vertical resolution.







IMG spectrum (WINCE, 970128 over Nebraska) and HITRAN database

Resolving absorption features in atmospheric windows enables detection of temperature inversions



Detection of inversions is critical for severe weather forecasting. Combined with improved low-level moisture depiction, key ingredients for night-time severe storm development can be monitored. 104

IASI detection of temperature inversion (black spectrum) VS clear ocean (red spectrum)





Ability to detect inversions disappears with broadband observations (> 3 cm-1)



Longwave window region



Longwave window region



Longwave window region


Longwave window region



Longwave window region



Longwave window region



Longwave window region

Twisted Ribbon formed by CO₂ spectrum: Tropopause inversion causes On-line & off-line patterns to cross













Inferring surface properties with AIRS high spectral resolution data Barren region detection if T1086 < T981 $T(981 \text{ cm}^{-1})$ - $T(1086 \text{ cm}^{-1})$

Barren vs Water/Vegetated



AIRS data from 14 June 2002



Intro to Lab on Split Window Moisture

Lecture - Labs in Bologna & Potenza May-June 2012

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Earth emitted spectra overlaid on Planck function envelopes



MODIS IR Spectral Bands

High resolution atmospheric absorption spectrum



First order estimation of SST correcting for low level moisture

Moisture attenuation in atmospheric windows varies linearly with optical depth.

$$\tau_{\lambda} = e = 1 - k_{\lambda} u$$

For same atmosphere, deviation of brightness temperature from surface temperature is a linear function of absorbing power. Thus moisture corrected SST can inferred by using split window measurements and extrapolating to zero k_{λ}

$$T_s = T_{bw1} + [k_{w1} / (k_{w2} - k_{w1})] [T_{bw1} - T_{bw2}]$$

Moisture content of atmosphere inferred from slope of linear relation.







In the IRW - A is off H2O line and B is on H2O line



Radiation is governed by Planck's Law

$$c_2 / \lambda T$$

B(λ ,T) = $c_1 / \{ \lambda^5 [e -1] \}$

In microwave region $c_2/\lambda T \ll 1$ so that $c_2/\lambda T$ $e = 1 + c_2/\lambda T + second order$

And classical Rayleigh Jeans radiation equation emerges

 $\mathbf{B}_{\lambda}(\mathbf{T}) \approx [\mathbf{c}_1 / \mathbf{c}_2] [\mathbf{T} / \lambda^4]$

Radiance is linear function of brightness temperature.

ISCCP-DX 199207-199306 Mean Annual



ISCCP-D1 1992 Mean Annual



Microwave Form of RTE

$$\frac{\text{ave Form of RTE}}{I^{\text{sfc}} = \varepsilon_{\lambda} B_{\lambda}(T_{s}) \tau_{\lambda}(p_{s}) + (1-\varepsilon_{\lambda}) \tau_{\lambda}(p_{s}) \int_{0}^{p_{s}} B_{\lambda}(T(p)) \frac{\partial \tau'_{\lambda}(p)}{\partial \ln p} d\ln p$$

$$I_{\lambda} = \varepsilon_{\lambda} B_{\lambda}(T_{s}) \tau_{\lambda}(p_{s}) + (1-\varepsilon_{\lambda}) \tau_{\lambda}(p_{s}) \int_{0}^{p_{s}} B_{\lambda}(T(p)) \frac{\partial \tau'_{\lambda}(p)}{\partial \ln p} d\ln p$$

$$+ \int_{p_{s}}^{0} B_{\lambda}(T(p)) \frac{\partial \tau_{\lambda}(p)}{\partial \ln p} d\ln p$$

$$\frac{\text{atm}}{f_{\lambda}(p)} d\ln p$$

$$\frac{d}{d} \ln p$$

$$\frac{d}{d} \ln p$$

In the microwave region $c_2/\lambda T \ll 1$, so the Planck radiance is linearly proportional to the brightness temperature

$$\mathsf{B}_{\lambda}(\mathsf{T}) \approx [\mathsf{c}_1 / \mathsf{c}_2] [\mathsf{T} / \lambda^4]$$

So

$$T_{b\lambda} = \epsilon_{\lambda} T_{s}(p_{s}) \tau_{\lambda}(p_{s}) + \int_{p_{s}}^{0} T(p) F_{\lambda}(p) \frac{\partial \tau_{\lambda}(p)}{\partial \ln p} d \ln p$$

where

$$F_{\lambda}(p) = \left\{ 1 + (1 - \varepsilon_{\lambda}) \left[\frac{\tau_{\lambda}(p_s)}{\tau_{\lambda}(p)} \right]^2 \right\}.$$

Transmittance

$$\tau(a,b) = \tau(b,a)$$

$$\tau(a,c) = \tau(a,b) * \tau(b,c)$$

Thus downwelling in terms of upwelling can be written

$$\tau'(p,ps) = \tau(ps,p) = \tau(ps,0) / \tau(p,0)$$

and

$$d\tau'(p,ps) = - d\tau(p,0) * \tau(ps,0) / [\tau(p,0)]^2$$

WAVELENGTH			FREQUENCY		WAVENUMBER
cm	μM	Å	Hz	GHz	cm⁻¹
10 ⁻⁵ Near Ultraviolet (I	0.1 UV)	1,000	3x10 ¹⁵		
4x10 ⁻⁵ Visible	0.4	4,000	7.5x10 ¹⁴		
7.5x10 ⁻⁵ Near Infrared (IR)	0.75)	7,500	4x10 ¹⁴		13,333
2x10 ⁻³ Far Infrared (IR)	20	2x10 ⁵	1.5x10 ¹³		500
0.1 Microwave (MW)	10 ³		3x10 ¹¹	300	10





Microwave spectral bands

- 23.8 GHz dirty window H2O absorption
- 31.4 GHz window
- 60 GHz O2 sounding
- 120 GHz O2 sounding
- 183 GHz H2O sounding







-270

-210

0·

-10-















 $Tb = \mathbf{\varepsilon} s T s \mathbf{\tau} m + \mathbf{\varepsilon} m T m + \mathbf{\varepsilon} m \mathbf{r} s \mathbf{\tau} m T m$

 $Tb = \varepsilon Ts (1-\sigma m) + \sigma m Tm + \sigma m (1-\varepsilon s) (1-\sigma m) Tm \quad using e^{-\sigma} = 1 - \sigma$

So temperature difference of low moist over ocean from clear sky over ocean is given by

 $\Delta Tb = - \varepsilon s \sigma m Ts + \sigma m Tm + \sigma m (1-\varepsilon s) (1-\sigma m) Tm$

For $\varepsilon_s \sim 0.5$ and $T_s \sim T_m$ this is always positive for $0 < \sigma_m < 1$



 $R = \varepsilon Bs (1-\sigma m) + \sigma m Bm$ using $e^{-\sigma} = 1 - \sigma$ and $\tau \sim 1-\sigma \sim 1-a$

So difference of low mist over ocean from clear sky over ocean is given by

 $\Delta R = - \varepsilon_s \sigma_m B_s + \sigma_m B_m$

For $\boldsymbol{\varepsilon}s \sim 1$

 $\Delta \mathbf{R} = -\boldsymbol{\sigma}_{m} \mathbf{B}_{s} + \boldsymbol{\sigma}_{m} \mathbf{B}_{m} = \boldsymbol{\sigma}_{m} [\mathbf{B}_{m} - \mathbf{B}_{s}]$

So if $[B_m - B_s] < 0$ then as σ_m increases ΔR becomes more negative







190

60N

SON

ΕQ

30S

60S-

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120E



120₩

AMSU 52.8

53.6

54.4 GHz

147

-210

205

-65

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AMSU 54.4

54.9

55.5 GHz


Spectral regions used for remote sensing of the earth atmosphere and surface from satellites. ε indicates emissivity, q denotes water vapour, and T represents temperature.